

Newton-Cartan Connections with Torsion

Tekin Dereli*

*Department of Physics, Koç University
34450 Sarıyer, İstanbul, Turkey*

Şahin Koçak[†] and Murat Limoncu[‡]

*Department of Mathematics, Anadolu University
26470 Eskişehir, Turkey*

February 7, 2008

Abstract

We re-formulate the notion of a Newton-Cartan manifold and clarify the compatibility conditions of a connection with torsion with the Newton-Cartan structure.

1 Introduction

According to Einstein's theory of general relativity, the space-time (M, g, ∇) is a four dimensional differentiable manifold M equipped with a type-(0,2), symmetric, non-degenerate, Lorentzian signed metric tensor g and the unique metric-compatible, torsion-free Levi-Civita connection ∇ . Elie Cartan [1] was first to notice that Newton's theory of gravitation can be formulated in a 4-dimensional space-time with a corank-1 degenerate metric [2], [3], [4]. The physical significance of Newton-Cartan space-times had been recognised after the 1960's [5], [6], [7]. It is well known that Cartan is also responsible for the concept of space-time torsion. In what follows we clarify the notion of Newton-Cartan manifolds and connections with torsion. In our main theorem (Th.2) we characterize n-dimensional Newton-Cartan space-times with torsion.

*E.mail: tdereli@ku.edu.tr

[†]E.mail: skocak@anadolu.edu.tr

[‡]E.mail: mlimoncu@anadolu.edu.tr

2 Newton-Cartan connections

Definition: 1 Let M be a smooth n -manifold, g and τ be (smooth) tensor fields of type $(0, 2)$ and $(0, 1)$, respectively, on M satisfying the following properties:

(G1) g is symmetric and degenerate with $\text{rank}(g) = n - 1$,

(G2) There exists a vector field $V \in \mathcal{X}(M)$ such that, $g(V, X) = 0$ for all $X \in \mathcal{X}(M)$ and $\tau(V) = 1$.

We call such a triplet (M, g, τ) a Newton-Cartan manifold.

Remark: The vector field V in (G2) which is called the time vector field is determined up to a scalar function: If $g_p(V_p, X_p) = 0$ and $g_p(W_p, X_p) = 0$ for all $X_p \in T_p M$ ($p \in M$) and for two vectors $V_p, W_p \in T_p M$, then V_p and W_p must be linearly dependent. Otherwise, V_p and W_p could be extended to a basis of $T_p M$, giving $\text{rank}(g) \leq n - 2$ and thus contradicting (G1).

Convention: Given any non-degenerate $(0, 2)$ -tensor \tilde{g} on a finite dimensional vector space V , a vector $X \in V$ determines a 1-form X^* , called the \tilde{g} -dual of X , by $X^*(Y) = \tilde{g}(X, Y)$ for any $Y \in V$. Likewise, a 1-form $\alpha \in V^*$ determines a vector α^* , called the \tilde{g} -dual of α , by $\tilde{g}(\alpha^*, Y) = \alpha(Y)$. We have $(X^*)^* = X$ and $(\alpha^*)^* = \alpha$. We will be using also the dual base (f^1, \dots, f^n) of V^* corresponding to a base (e_1, \dots, e_n) of V , given by $f^i(e_j) = \delta_j^i$, which is to be distinguished from the \tilde{g} -dual base (e_1^*, \dots, e_n^*) .

We will consider a fixed vector field V with $g(V, X) = 0$ and $\tau(V) = 1$ also as part of a Newton-Cartan manifold. g restricted to the kernel of τ (more precisely to $\ker \tau \times \ker \tau$) is non-degenerate and one easily obtains:

Proposition:1 The tensor field $\bar{g}_\epsilon = g + \epsilon \tau \otimes \tau$, where ϵ is a non-zero real number, is symmetric and non-degenerate. Thus it is a semi-Riemannian metric on M . The dual τ^* of τ with respect to \bar{g}_ϵ is $\frac{1}{\epsilon}V$, i.e. $\tau(X) = \bar{g}_\epsilon(\frac{1}{\epsilon}V, X)$.

Remark: In the following, unless necessary, we will drop the subscript ϵ and understand that a fixed ϵ is chosen. We will stress the role of ϵ in Theorems 1 and 2.

Let us define $\bar{h}(\alpha, \beta) = \bar{g}(\alpha^*, \beta^*)$, where α and β are 1-forms and α^*, β^* their \bar{g} -duals. The $(2, 0)$ -tensor field \bar{h} is symmetric and non-degenerate.

Proposition: 2 The $(2, 0)$ -tensor field $h = \bar{h} - \frac{1}{\epsilon}V \otimes V$ is symmetric and degenerate.

Proof: h is obviously symmetric. For degeneracy, let θ be any 1-form field. Then

$$\begin{aligned}
h(\tau, \theta) &= \bar{h}(\tau, \theta) - \frac{1}{\epsilon} V(\tau) V(\theta) \\
&= \bar{g}(\tau^*, \theta^*) - \frac{1}{\epsilon} \tau(V) \theta(V) \\
&= \bar{g}\left(\frac{1}{\epsilon} V, \theta^*\right) - \frac{1}{\epsilon} \theta(V) \\
&= \frac{1}{\epsilon} (\bar{g}(V, \theta^*) - \theta(V)) = 0.
\end{aligned}$$

Proposition:3 The following relations between \bar{h} and \bar{g} (resp. h and g) hold:

$$\begin{aligned}
c_1^2(\bar{h} \otimes \bar{g}) &= \delta, \\
c_1^2(h \otimes g) &= \delta - V \otimes \tau
\end{aligned}$$

where c_1^2 denotes the $(2, 1)$ -contraction of the $(2, 2)$ -tensor $\bar{h} \otimes \bar{g}$ (resp. $h \otimes g$) and δ is the $(1, 1)$ -tensor $\delta(\alpha, X) = \alpha(X)$. In a local coordinate chart (x^1, \dots, x^n) we have in terms of components

$$\begin{aligned}
\sum_k \bar{h}^{ik} \bar{g}_{kj} &= \delta_j^i, \\
\sum_k h^{ik} g_{kj} &= \delta_j^i - V^i \tau_j.
\end{aligned}$$

These equalities are also true for any non-coordinate base $\{e_i\}$ and its dual base $\{f^j\}$ with $f^j(e_i) = \delta_i^j$.

Proof: Let $(\partial_1, \dots, \partial_n)$ denote the coordinate base, (dx^1, \dots, dx^n) the dual base. Since $\bar{g}(\partial_i, \partial_j) = \bar{g}_{ij}$ and $\bar{h}(dx^i, dx^j) = \bar{h}^{ij}$ we have by definition $\sum_k \bar{h}^{ik} \bar{g}_{kj} = \delta_j^i$. Inserting in this $\bar{h}^{ik} = h^{ik} + \frac{1}{\epsilon} \tau^i \tau^k$ and $\bar{g}_{kj} = g_{kj} + \epsilon V_k V_j$ and taking into account $\tau(V) = 1$, $h(dx^i, \tau) = 0$, $g(V, \partial_j) = 0$ we get $\sum_k h^{ik} g_{kj} = \delta_j^i - \tau^i V_j$.

Having defined a semi-Riemannian manifold (M, \bar{g}) associated with the Newton-Cartan manifold (M, g, τ) , we have the standard theory of connections at our disposal. Imposing certain conditions on connections on M , we will single out *Newton-Cartan* connections on (M, g, τ) compatible with g and τ in the following sense:

Definition: 2 Let ∇_X denote the covariant derivative of any linear connection \mathcal{D} on the manifold M . We will call the connection \mathcal{D} a *Newton-Cartan* connection iff

$$\nabla_X \tau = 0 \quad , \quad \nabla_X g = 0 \quad \forall X \in \mathcal{X}(M)$$

Proposition: 4 Let ∇_X denote the covariant derivative of a *Newton-Cartan* connection on M . Then $\nabla_X V = 0$ and $\nabla_X h = 0$, where V is the time vector field and h the (2,0)-field associated with g (s. Prop 2). Moreover, $\nabla_X \bar{h} = 0$. Proof: First we note that $\nabla_X \bar{g} = \nabla_X g + \epsilon(\nabla_X \tau \otimes \tau + \tau \otimes \nabla_X \tau) = 0$. Now, from Prop.1, $\epsilon\tau(Y) = \bar{g}(V, Y)$ for any $Y \in \mathcal{X}(M)$. Then

$$\begin{aligned} \nabla_X(\epsilon\tau(Y)) &= \nabla_X(\bar{g}(V, Y)) \\ \epsilon(\nabla_X \tau)(Y) + \epsilon\tau(\nabla_X Y) &= (\nabla_X \bar{g})(V, Y) + \bar{g}(\nabla_X V, Y) + \bar{g}(V, \nabla_X Y) \\ \epsilon\tau(\nabla_X Y) &= \bar{g}(\nabla_X V, Y) + \bar{g}(V, \nabla_X Y) \end{aligned}$$

by $\nabla_X \tau = 0$ and $\nabla_X \bar{g} = 0$. Since $\epsilon\tau(\nabla_X Y) = \bar{g}(V, \nabla_X Y)$ we get $\bar{g}(\nabla_X V, Y) = 0$, hence $\nabla_X V = 0$.

Now we consider $h = \bar{h} - \frac{1}{\epsilon} V \otimes V$

$$\begin{aligned} \nabla_X h &= \nabla_X \bar{h} - \frac{1}{\epsilon} V \otimes \nabla_X V - \frac{1}{\epsilon} \nabla_X V \otimes V \\ \nabla_X h &= \nabla_X \bar{h} \end{aligned}$$

So it is enough to show $\nabla_X \bar{h} = 0$. Let α, β be 1-form fields and α^*, β^* their \bar{g} -dual vector fields. Since by definition $\bar{h}(\alpha, \beta) = \bar{g}(\alpha^*, \beta^*)$,

$$(\nabla_X \bar{h})(\alpha, \beta) = \nabla_X(\bar{h}(\alpha, \beta)) - \bar{h}(\nabla_X \alpha, \beta) - \bar{h}(\alpha, \nabla_X \beta).$$

On the other hand we have $(\nabla_X \alpha)^* = \nabla_X(\alpha^*)$:

$$\begin{aligned} \bar{g}(\alpha^*, U) &= \alpha(U) \quad U \in \mathcal{X}(M) \\ \nabla_X(\bar{g}(\alpha^*, U)) &= \nabla_X(\alpha(U)) \\ \bar{g}(\nabla_X \alpha^*, U) + \bar{g}(\alpha^*, \nabla_X U) &= (\nabla_X \alpha)(U) + \alpha(\nabla_X U) \\ \bar{g}(\nabla_X \alpha^*, U) &= (\nabla_X \alpha)(U) \end{aligned}$$

because $\bar{g}(\alpha^*, \nabla_X U) = \alpha(\nabla_X U)$. This means $(\nabla_X \alpha)^* = \nabla_X \alpha^*$. Inserting these correspondences above and from $\nabla_X \bar{g} = 0$ we obtain

$$(\nabla_X \bar{h})(\alpha, \beta) = \nabla_X(\bar{g}(\alpha^*, \beta^*)) - \bar{g}(\nabla_X \alpha^*, \beta^*) - \bar{g}(\alpha^*, \nabla_X \beta^*) = 0.$$

Remark: The converse of this proposition is also true and a kind of duality emerges. One could define a *Newton-Cartan* manifold also as a triplet (M, h, V) .

The following theorem clarifies the existence and uniqueness of torsion-free *Newton-Cartan* connections:

Theorem:1 Let (M, g, τ) be a *Newton-Cartan* manifold.

i) If \mathcal{D} is a *Newton-Cartan* connection without torsion, then \mathcal{D} is the Levi-Civita connection on (M, \bar{g}_ϵ) for all $\epsilon (\neq 0)$ and $L_V g = 0, d\tau = 0$ where L_V denotes the Lie derivative, V being the time vector field ($\tau(V) = 1$).

ii) Let \mathcal{D} denote the Levi-Civita connection on (M, \bar{g}_ϵ) for a fixed $\epsilon (\neq 0)$. Assume, $L_V g = 0$ and $d\tau = 0$. Then, \mathcal{D} is a torsion-free *Newton-Cartan* connection on (M, g, τ) .

Remark: Theorem 1 shows that, there can't be any torsion-free *Newton-Cartan* connection, unless $L_V g = 0$ and $d\tau = 0$. In case of $L_V g = 0$ and $d\tau = 0$, there is a unique *Newton-Cartan* connection without torsion; it is the Levi-Civita connection for any \bar{g}_ϵ (especially, Levi-Civita connections for \bar{g}_ϵ coincide for all $\epsilon \neq 0$).

Proof (of theorem 1):

i) The first assertion (that $\nabla_X \bar{g}_\epsilon = 0$) is trivial. So we show $d\tau = 0$ and $L_V g = 0$.

$$\begin{aligned} (\nabla_X \tau)(Y) &= \nabla_X (\tau(Y)) - \tau(\nabla_X Y) = 0 \\ (\nabla_Y \tau)(X) &= \nabla_Y (\tau(X)) - \tau(\nabla_Y X) = 0 \end{aligned}$$

$$\begin{aligned} X\tau(Y) - Y\tau(X) &= \tau(\nabla_X Y - \nabla_Y X) = \tau([X, Y]) \\ (d\tau)(X, Y) &= 0 \quad d\tau = 0. \end{aligned}$$

Now we will show $L_V g = 0$:

$$\begin{aligned} (\nabla_X g)(Y, Z) &= \nabla_X (g(Y, Z)) - g(\nabla_X Y, Z) - g(Y, \nabla_X Z) = 0 \\ &= Xg(Y, Z) - g([X, Y] + \nabla_Y X, Z) - g(Y, [X, Z] + \nabla_Z X) = 0 \\ &= Xg(Y, Z) - g(L_X Y, Z) - g(Y, L_X Z) - g(\nabla_Y X, Z) - g(Y, \nabla_Z X) = 0 \end{aligned}$$

Take $X = V$:

$$Vg(Y, Z) - g(L_V Y, Z) - g(Y, L_V Z) - g(\nabla_Y V, Z) - g(Y, \nabla_Z V) = 0$$

By Prop.4, $\nabla_Y V = 0$ $\nabla_Z V = 0$ and thus

$$Vg(Y, Z) - g(L_V Y, Z) - g(Y, L_V Z) = 0 \implies (L_V g)(Y, Z) = 0 (L_V g = 0).$$

ii) It is enough to show $\nabla_X \tau = 0$ as $\nabla_X g = 0$ is then a consequence of $\bar{g}_\epsilon = g + \epsilon \tau \otimes \tau$ and $\nabla_X \bar{g}_\epsilon = 0$. (We will write below \bar{g} for \bar{g}_ϵ again)

Let $Y \in \mathcal{X}(M)$. We get

$$\begin{aligned} (\nabla_X \tau)(Y) &= \nabla_X(\tau(Y)) - \tau(\nabla_X Y) \\ (\nabla_X \tau)(Y) &= X\left(\bar{g}\left(\frac{1}{\epsilon}V, Y\right)\right) - \bar{g}\left(\frac{1}{\epsilon}V, \nabla_X Y\right) \\ \epsilon(\nabla_X \tau)(Y) &= X(\bar{g}(V, Y)) - \bar{g}(V, \nabla_X Y) \end{aligned}$$

since $\frac{1}{\epsilon}V$ is \bar{g} -dual of τ . From the Koszul formula

$$\begin{aligned} 2\bar{g}(\nabla_X Y, V) &= X\bar{g}(Y, V) + Y\bar{g}(V, X) - V\bar{g}(X, Y) \\ &\quad - \bar{g}(X, [Y, V]) + \bar{g}(Y, [V, X]) + \bar{g}(V, [X, Y]) \end{aligned}$$

we can write

$$\begin{aligned} \epsilon(\nabla_X \tau)(Y) &= \frac{1}{2}\{X\bar{g}(V, Y) - Y\bar{g}(V, X) - \bar{g}(V, [X, Y])\} \\ &\quad + \frac{1}{2}\{V\bar{g}(X, Y) - \bar{g}(X, [V, Y]) - \bar{g}(Y, [V, X])\}. \end{aligned} \quad (1)$$

The first term vanishes because of $d\tau = 0$:

$$\begin{aligned} (d\tau)(X, Y) &= X(\tau(Y)) - Y(\tau(X)) - \tau([X, Y]) = 0 \\ &= X\left(\bar{g}\left(\frac{1}{\epsilon}V, Y\right)\right) - Y\left(\bar{g}\left(\frac{1}{\epsilon}V, X\right)\right) - \bar{g}\left(\frac{1}{\epsilon}V, [X, Y]\right) = 0. \end{aligned}$$

The second term in (1) vanishes because of $L_V \bar{g} = 0$: (we will instantly show that $L_V \bar{g} = 0$)

$$\begin{aligned} (L_V \bar{g})(X, Y) &= V\bar{g}(X, Y) - \bar{g}(L_V X, Y) - \bar{g}(X, L_V Y) = 0 \\ &= V\bar{g}(X, Y) - \bar{g}(X, [V, Y]) - \bar{g}(Y, [V, X]) = 0 \end{aligned}$$

Now the missing last part of the proof: $L_V \bar{g} = 0$. First we show that $L_V \tau = 0$:

$$\begin{aligned} (L_V \tau)(X) &= L_V(\tau(X)) - \tau(L_V X) \\ &= V(\tau(X)) - \tau([V, X]) \\ &= X(\tau(V)) \quad (d\tau = 0) \\ &= 0 \quad (\tau(V) = 1) \end{aligned}$$

Now,

$$\begin{aligned} L_V \bar{g} &= L_V g + \epsilon L_V(\tau \otimes \tau) \\ L_V \bar{g} &= L_V g + \epsilon L_V \tau \otimes \tau + \epsilon \tau \otimes L_V \tau \\ L_V \bar{g} &= L_V g = 0 \end{aligned}$$

by assumption of the theorem.

We generalize to the cases with torsion and prove our main theorem below:

Theorem:2 Let (M, g, τ) be a *Newton-Cartan* manifold

i) If \mathcal{D} is a *Newton-Cartan* connection with torsion T ($T(X, Y) = \nabla_X Y - \nabla_Y X - [X, Y]$) then \mathcal{D} is \bar{g}_ϵ compatible for all $\epsilon (\neq 0)$ (i.e. $\nabla_X \bar{g}_\epsilon = 0$) and

$$\begin{aligned} (L_V g)(X, Y) &= g(T(V, X), Y) + g(X, T(V, Y)) \\ (d\tau)(X, Y) &= \tau(T(X, Y)) \end{aligned} \quad (2)$$

for all $X, Y, Z \in \mathcal{X}(M)$.

ii) Let \mathcal{D} denote the connection on (M, \bar{g}_ϵ) , compatible with \bar{g}_ϵ and with torsion T (for a fixed $\epsilon \neq 0$). Assume

$$\begin{aligned} (L_V g)(X, Y) &= g(T(V, X), Y) + g(X, T(V, Y)) \\ (d\tau)(X, Y) &= \tau(T(X, Y)) \end{aligned}$$

for all $X, Y, Z \in \mathcal{X}(M)$. Then \mathcal{D} is a *Newton-Cartan* connection on (M, g, τ) with torsion T .

Remark: We call the conditions (2) *compatibility conditions* of a torsion with the *Newton-Cartan* structure. Theorem 2 shows that, a *Newton-Cartan* connection with torsion T can exist, only if the torsion T is compatible with the *Newton-Cartan* structure. On the other hand, if T is compatible with the *Newton-Cartan* structure, then there exists a unique *Newton-Cartan* connection with torsion T ; it is the connection compatible with any \bar{g}_ϵ and having torsion T (especially, connections compatible with \bar{g}_ϵ and having torsion T coincide for all $\epsilon \neq 0$).

Proof (of theorem 2)

i) The first assertion (that $\nabla_X \bar{g}_\epsilon = 0$) is again trivial. So we show the other two identities. First we show $(d\tau)(X, Y) = \tau(T(X, Y))$:

$$\begin{aligned} (\nabla_X \tau)(Y) &= \nabla_X (\tau(Y)) - \tau(\nabla_X Y) = 0 \\ (\nabla_Y \tau)(X) &= \nabla_Y (\tau(X)) - \tau(\nabla_Y X) = 0 \end{aligned}$$

$$\begin{aligned} X\tau(Y) - Y\tau(X) &= \tau(\nabla_X Y - \nabla_Y X) = \tau(T(X, Y) + [X, Y]) \\ (d\tau)(X, Y) &= X\tau(Y) - Y\tau(X) - \tau([X, Y]) = \tau(T(X, Y)). \end{aligned}$$

Now we show the identity about $L_V g$:

$$(\nabla_X g)(Y, Z) = \nabla_X (g(Y, Z)) - g(\nabla_X Y, Z) - g(Y, \nabla_X Z) = 0$$

$$Xg(Y, Z) - g(T(X, Y) + [X, Y] + \nabla_Y X, Z) - g(Y, T(X, Z) + [X, Z] + \nabla_Z X) = 0$$

$$\begin{aligned} Xg(Y, Z) - g(L_X Y, Z) - g(Y, L_X Z) \\ - g(T(X, Y) + \nabla_Y X, Z) - g(Y, T(X, Z) + \nabla_Z X) = 0 \end{aligned}$$

Take $X = V$:

$$Vg(Y, Z) - g(L_V Y, Z) - g(Y, L_V Z) - g(T(V, Y) + \nabla_Y V, Z) - g(Y, T(V, Z) + \nabla_Z V) = 0$$

By Prop.4, $\nabla_Y V = 0$ $\nabla_Z V = 0$ and thus

$$\begin{aligned} Vg(Y, Z) - g(L_V Y, Z) - g(Y, L_V Z) &= g(T(V, Y), Z) + g(Y, T(V, Z)) \\ (L_V g)(Y, Z) &= g(T(V, Y), Z) + g(Y, T(V, Z)). \end{aligned}$$

ii) Again, it is enough to show $\nabla_X \tau = 0$ as $\nabla_X g = 0$ follows from $\bar{g} = g + \epsilon \tau \otimes \tau$ and $\nabla_X \bar{g} = 0$.

The equality for $\nabla_X \tau$,

$$\epsilon(\nabla_X \tau)(Y) = X(\bar{g}(V, Y)) - \bar{g}(V, \nabla_X Y)$$

is still true. But now we have the Koszul formula with torsion:

$$\begin{aligned} 2\bar{g}(\nabla_X Y, V) &= X\bar{g}(Y, V) + Y\bar{g}(V, X) - V\bar{g}(X, Y) \\ &\quad - \bar{g}(X, [Y, V]) + \bar{g}(Y, [V, X]) + \bar{g}(V, [X, Y]) \\ &\quad - \bar{g}(X, T(Y, V)) + \bar{g}(Y, T(V, X)) + \bar{g}(V, T(X, Y)) \end{aligned}$$

where $T(X, Y) = \nabla_X Y - \nabla_Y X - [X, Y]$.

Inserting this into the equality for $\nabla_X \tau$ we get

$$\begin{aligned} 2\epsilon(\nabla_X \tau)(Y) &= \{X\bar{g}(V, Y) - Y\bar{g}(V, X) - \bar{g}(V, [X, Y])\} \\ &\quad + \{V\bar{g}(X, Y) - \bar{g}(X, [V, Y]) - \bar{g}(Y, [V, X])\} \\ &\quad + \{\bar{g}(X, T(Y, V)) - \bar{g}(Y, T(V, X)) - \bar{g}(V, T(X, Y))\}. \end{aligned} \tag{3}$$

We first simplify the first term in (3):

By assumption $(d\tau)(X, Y) = \tau(T(X, Y))$

$$\begin{aligned} X(\tau(Y)) - Y(\tau(X)) - \tau([X, Y]) &= \tau(T(X, Y)) \\ X\left(\bar{g}\left(\frac{1}{\epsilon}V, Y\right)\right) - Y\left(\bar{g}\left(\frac{1}{\epsilon}V, X\right)\right) - \bar{g}\left(\frac{1}{\epsilon}V, [X, Y]\right) &= \bar{g}\left(\frac{1}{\epsilon}V, T(X, Y)\right) \end{aligned}$$

Thus, the first term in (3) equals $\bar{g}(V, T(X, Y))$.

To simplify the second term in (3) we use as before,

$$\begin{aligned}(L_V \bar{g})(X, Y) &= V\bar{g}(X, Y) - \bar{g}(L_V X, Y) - \bar{g}(X, L_V Y) \\ &= V\bar{g}(X, Y) - \bar{g}(X, [V, Y]) - \bar{g}(Y, [V, X])\end{aligned}$$

Thus, the second term in (3) equals $(L_V \bar{g})(X, Y)$. Let us compute $L_V \bar{g}$: Since $\bar{g} = g + \epsilon\tau \otimes \tau$

$$L_V \bar{g} = L_V g + \epsilon L_V \tau \otimes \tau + \epsilon\tau \otimes L_V \tau$$

so we need $L_V \tau$:

$$\begin{aligned}(L_V \tau)(X) &= L_V(\tau(X)) - \tau(L_V X) \\ &= V(\tau(X)) - \tau([V, X]) \\ &= X(\tau(V)) + \tau(T(V, X)) \quad (d\tau)(V, X) = \tau(T(V, X)) \\ &= \tau(T(V, X)) \quad (\tau(V) = 1)\end{aligned}$$

Inserting this in the formula for $L_V \bar{g}$, we obtain

$$\begin{aligned}(L_V \bar{g})(X, Y) &= (L_V g)(X, Y) + \epsilon(L_V \tau)(X)\tau(Y) + \epsilon\tau(X)(L_V \tau)(Y) \\ &= g(T(V, X), Y) + g(X, T(V, Y)) \\ &\quad + \epsilon\tau(T(V, X))\tau(Y) + \epsilon\tau(X)\tau(T(V, Y)) \\ &= \bar{g}(T(V, X), Y) + \bar{g}(X, T(V, Y))\end{aligned}$$

If we now replace the first term in (3) by $\bar{g}(V, T(X, Y))$ and second term in (3) by $\bar{g}(T(V, X), Y) + \bar{g}(X, T(V, Y))$ we get $\nabla_X \tau = 0$ (by anti-symmetry of T).

Remark: There are two cases where the conditions of Theorem 2 simplify. Namely,

- i. If $T(X, Y) \in \ker \tau$, then the condition $(d\tau)(X, Y) = \tau(T(X, Y))$ reduces to $d\tau = 0$.
- ii. If $T(X, Y) = d\tau(X, Y)V$ then the condition $(d\tau)(X, Y) = \tau(T(X, Y))$ is satisfied. Because, $\tau(T(X, Y)) = \tau(d\tau(X, Y)V) = d\tau(X, Y)\tau(V) = d\tau(X, Y)$ by $\tau(V) = 1$. In this case, the condition $(L_V g)(X, Y) = g(T(V, X), Y) + g(X, T(V, Y))$ reduces to $L_V g = 0$ because $g(T(V, X), Y) = g(d\tau(V, X)V, Y) = d\tau(V, X)g(V, Y) = 0$ and similarly $g(X, T(V, Y)) = 0$.

Emerging from the considerations in Theorem 1 and Theorem 2 (and remarks following them), there holds the following characterization of *Newton-Cartan* connections, which is interesting by itself.

Proposition 5: A connection \mathcal{D} (with or without torsion) on a *Newton-Cartan* manifold (M, g, τ) is a *Newton-Cartan* connection, if and only if it is compatible with all metrics $\bar{g}_\epsilon = g + \epsilon\tau \otimes \tau$ ($\epsilon \neq 0$).

(i.e. $\nabla_X g = 0$ and $\nabla_X \tau = 0$ for all $X \in \mathcal{X}(M)$) $\iff \nabla_X \bar{g}_\epsilon = 0$ for all $\epsilon \neq 0$ and $X \in \mathcal{X}(M)$)

Proof: It is obvious that $\nabla_X g = 0$ and $\nabla_X \tau = 0$ imply $\nabla_X \bar{g}_\epsilon = 0$. Now assume $\nabla_X \bar{g}_\epsilon = 0$ for all $\epsilon \neq 0$. Indeed, it will suffice to assume $\nabla_X \bar{g}_{\epsilon_1} = 0$ and $\nabla_X \bar{g}_{\epsilon_2} = 0$ for two specific and distinct $\epsilon_1 \neq \epsilon_2$ (e.g. ± 1). From $\nabla_X g + \epsilon_1 \nabla_X (\tau \otimes \tau) = 0$ and $\nabla_X g + \epsilon_2 \nabla_X (\tau \otimes \tau) = 0$ we get $\nabla_X g = 0$ and $\nabla_X (\tau \otimes \tau) = 0$. By the product rule

$$\begin{aligned} (\nabla_X (\tau \otimes \tau))(Y, Z) &= \nabla_X ((\tau \otimes \tau)(Y, Z)) - (\tau \otimes \tau)(\nabla_X Y, Z) - (\tau \otimes \tau)(Y, \nabla_X Z) = 0 \\ &= X(\tau(Y)\tau(Z)) - \tau(\nabla_X Y)\tau(Z) - \tau(Y)\tau(\nabla_X Z) = 0 \end{aligned}$$

for all $X, Y, Z \in \mathcal{X}(M)$. Taking $Y = Z = V$ we find $\tau(\nabla_X V) = 0$ by $\tau(V) = 1$. If we now take in the above identity only $Z = V$ and use $\tau(\nabla_X V) = 0$ we get

$$X(\tau(Y)) - \tau(\nabla_X Y) = 0$$

which means $(\nabla_X \tau)(Y) = 0$, i.e. $\nabla_X \tau = 0$.

3 Conclusion

We discuss here the relation between n -dimensional pseudo-Riemannian manifolds with parallel vector fields and Newton-Cartan structures. The latter consist of two parallel tensors, namely, a type- $(0, 2)$ symmetric tensor g that is degenerate of corank- $(n-1)$ and a non-vanishing 1-form τ such that g is non-degenerate on $\ker(\tau)$ and a compatible connection ∇ . From this data, it is possible to construct a 1-parameter family of pseudo-Riemannian metrics $\bar{g}_\epsilon = g + \epsilon\tau \otimes \tau$ whose Levi-Civita connection coincide for any ϵ with the (torsion-free) Newton-Cartan connection and the 1-form τ and its dual vector V are parallel. Conversely, from a pseudo-Riemannian metric and a parallel vector field V one may reconstruct the Newton-Cartan structure making use of the dual 1-form τ . We give formulas that characterize a metric g and a parallel vector field V that are compatible with a Newton-Cartan connection with torsion. Most of these results were previously announced in an unpublished thesis [8].

4 Acknowledgements

TD thanks Professors J. Ehlers and R. W. Tucker for discussions that motivated this work.

References

- [1] E. Cartan, *Sur les variétés à connexion affine et la théorie de la relativité généralisée*, Ann. Éc. Norm., **XL**, (1922) 325 , *ibid*, **XLI** (1924)1
- [2] K. Friedrichs, *Eine invariante Formulierung des Newtonschen Gravitationsgesetzes und des Grenzüberganges vom Einsteinschen zum Newtonschen Gesetz*, Math. Ann. **98** (1927) 566
- [3] H. D. Dombrowski, K. Horneffer, *Die Differentialgeometrie des Galileischen Relativitätsprinzips*, Math. Zeitschr. **86** (1964) 291
- [4] M. Crampin, *On differentiable manifolds with degenerate metrics*, Proc. Camb. Phil. Soc. **64** (1968) 307
- [5] P. Havas, *Four-dimensional formulations of Newton mechanics and their relation to the special and the general theory of relativity*, Rev. Mod. Phys. **36**, (1964) 938
- [6] A. Trautman, *Comparison of Newtonian and relativistic theories of space-time*, in **Perspectives in Geometry and Relativity** Edited by B. Hoffmann (Indiana University Press, 1966) p. 413
- [7] J. Christian, *Exactly soluble sector of quantum gravity*, Phys. Rev. **D56** (1997) 4844
- [8] M. Limoncu, *Differential geometry of Newton-Cartan theories of gravitation*, unpublished M. S. thesis (in Turkish) submitted to Anadolu University, January 2001.